

# Physics 1C

Lecture 28B

# Essential Concept of last lecture

- Light is emitted and absorbed in form of quantas of energy that are multiple of a minimal energy quantum.
  - Emission: e.g. black body radiation
  - Absorption: e.g. photoelectric effect
- This minimal quantum of energy depends of the frequency of the light as follows:

$$E_0 = hf \Rightarrow E_n = nhf$$

- $h=6.626 \times 10^{-34}$  Joule x seconds is a constant of nature called Planck's constant.

# The Compton Effect

- In 1923, Arthur Compton (U of Chicago) directed a beam of x-rays toward a block of graphite
- He detected the scattered x-rays had a slightly longer wavelength than the incident x-rays.
- This means the scattered photons had less energy than the incident photons.
- The amount of energy the scattered photons lost depended on the angle at which the x-rays were scattered.
- This change in wavelength is called the **Compton shift** (1927 Nobel Prize).

# The Compton Effect

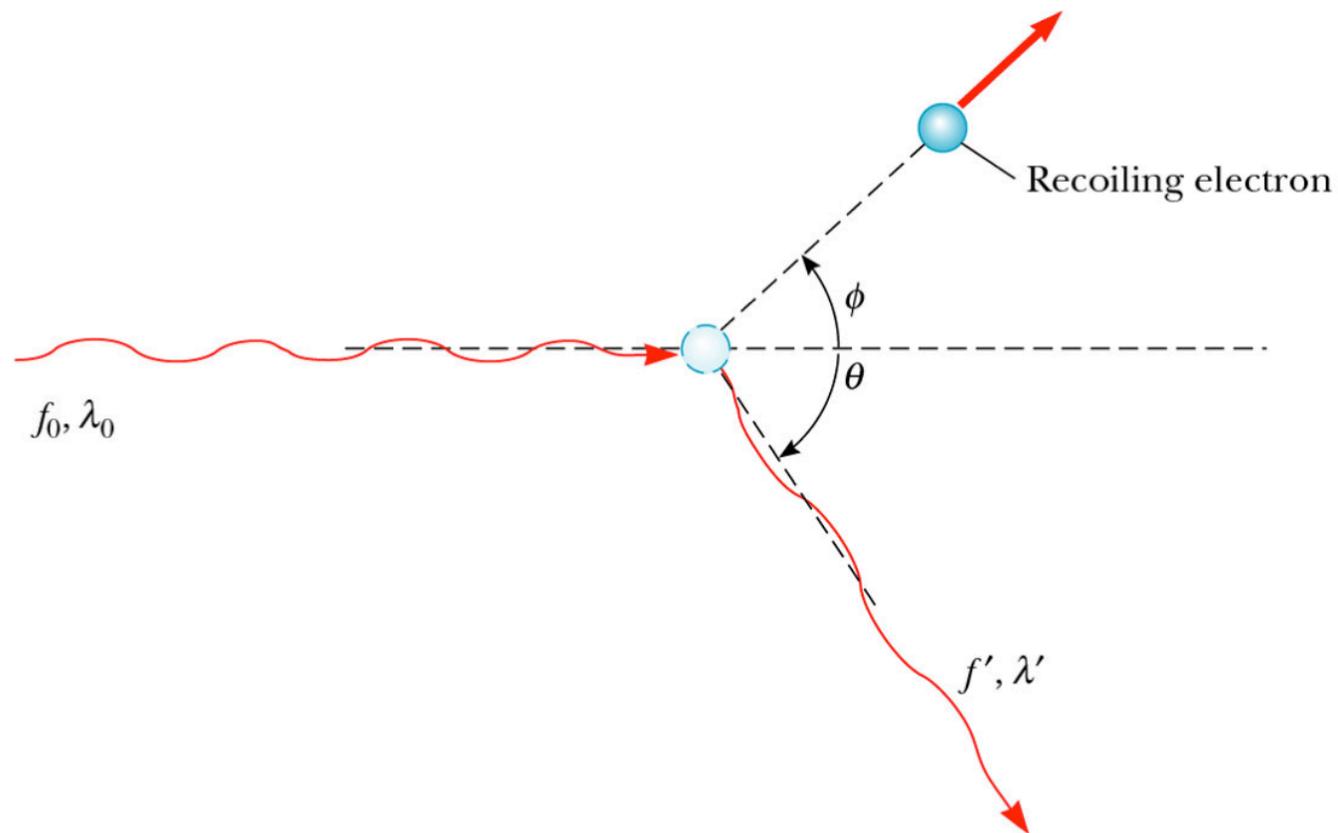
To calculate the shift in wavelength, Compton assumed that the photons act like other particles in collisions.

In the collisions, energy,  $hf$ , and momentum,  $hf/c$ , were conserved.

The energy of the incoming photon was:

$$E_o = hf_o = h \frac{c}{\lambda_o}$$

After it collides it scatters by an angle  $\theta$ .



# The Compton Effect

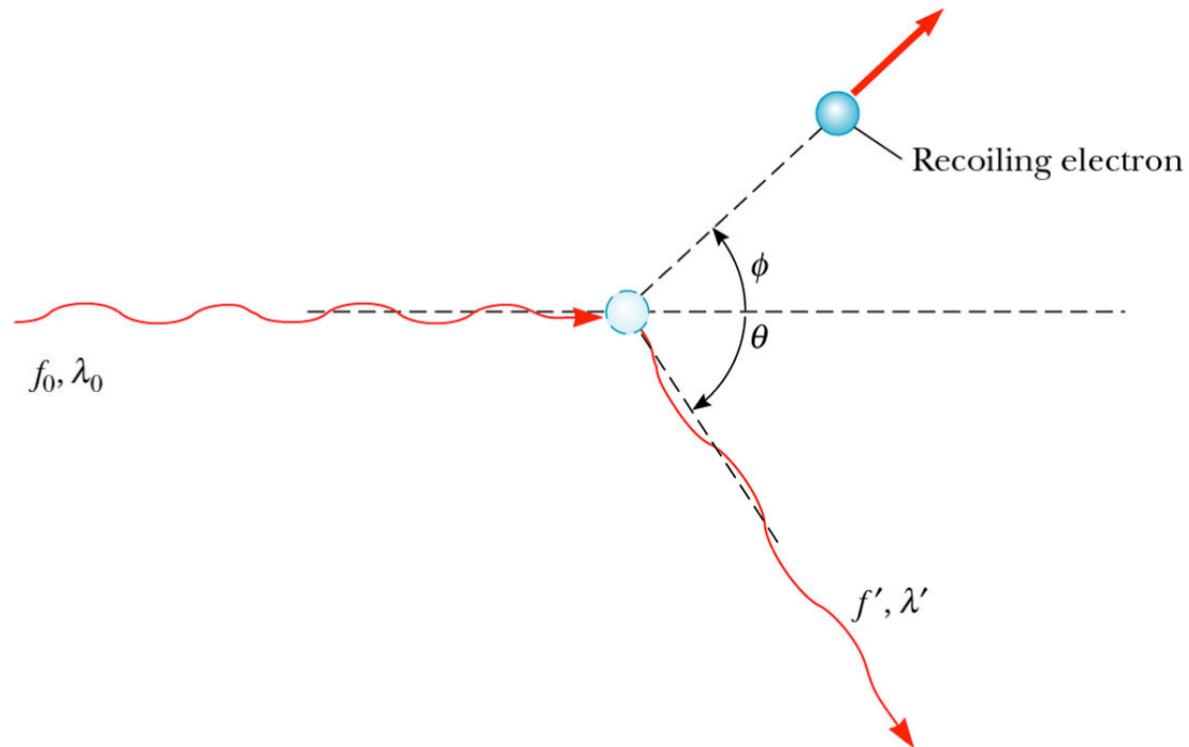
After the collision the photon has an energy:

$$E_f = hf = h \frac{c}{\lambda}$$

The wavelength shift becomes:  $\lambda' - \lambda_0 = \frac{h}{m_e c} (1 - \cos \theta)$

where the **Compton wavelength** for the electron is a constant:

$$\frac{h}{m_e c} = 0.00243 \text{ nm}$$



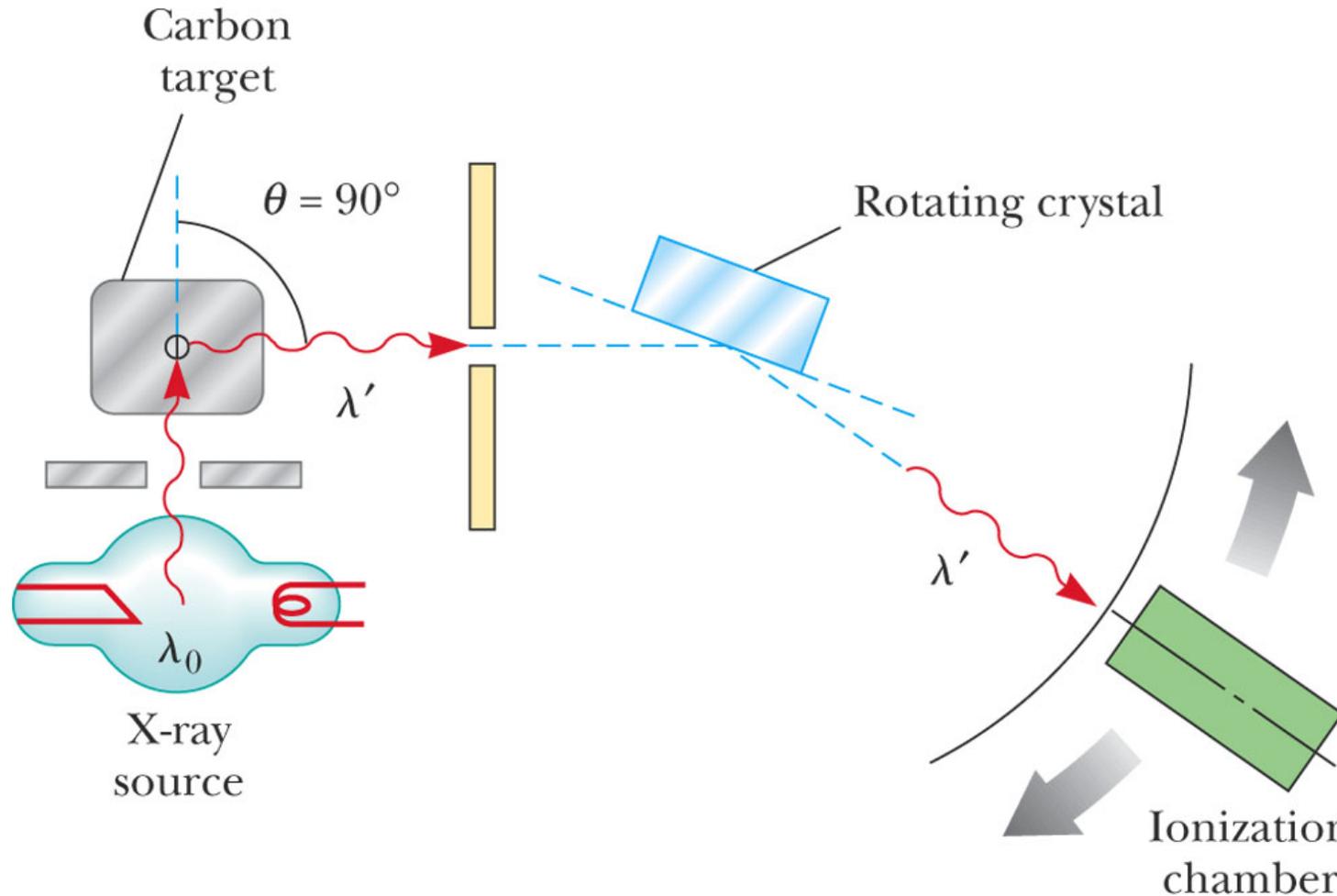
# The Compton Effect

Schematic diagram of Compton's apparatus:

An x-ray tube produces radiation of

$\lambda_0 = 0.0707\text{nm}$  that strikes a carbon target.

Angle  $\theta$  is varied by moving the x-ray source.



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The x-ray spectrometer includes a crystal that reflects x-rays and an ionization chamber that measures  $I$ .

# The Compton Effect

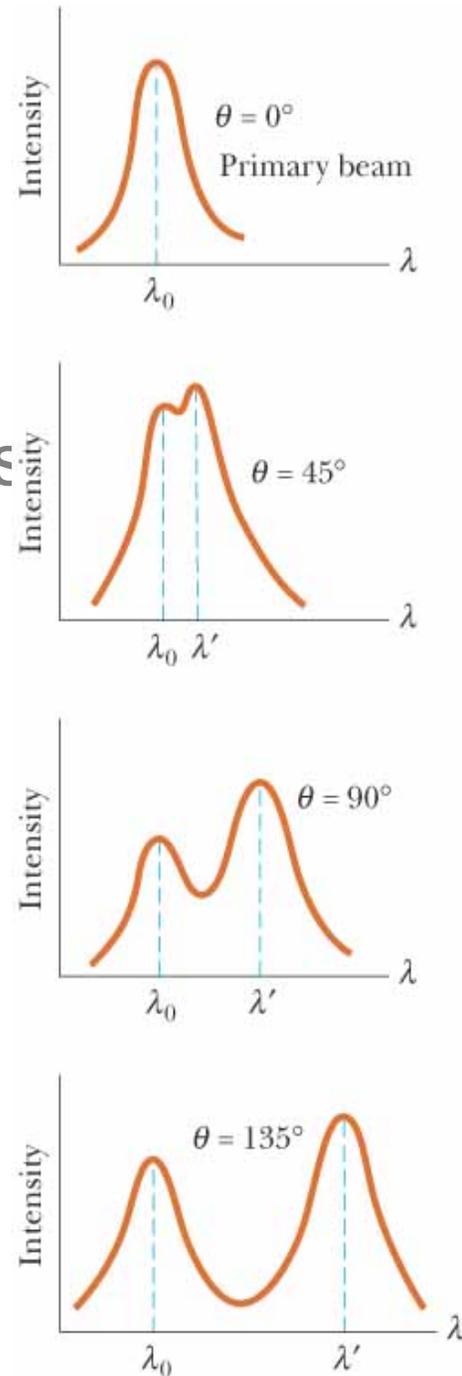
The wavelength of the scattered x-rays can be determined from the angle at which they were reflected from the crystal with maximum intensity (x-ray diffraction).

The graphs show the spectra of scattered x-rays for various angles  $\theta$ .

The shifted peak at  $\lambda'$  is caused by the scattering of free electrons in the target:

$$\lambda' - \lambda_0 = \frac{h}{m_e c} (1 - \cos \theta)$$

The unshifted wavelength,  $\lambda_0$ , is due to x-rays scattered from the electrons that are tightly bound to the target atoms.



# Wave Particle Duality

- Experiments such as the photoelectric effect, the Compton effect, and double-slit interference helped to describe light's true nature: a wave and a particle.
- Light has a dual nature, exhibiting both wave and particle characteristics.
- This means a photon can have energy given by:

$$E = hf = h \frac{c}{\lambda}$$

- And that photons can have a momentum given by:

$$p = \frac{E}{c} = \frac{1}{c} \left( h \frac{c}{\lambda} \right) = \frac{h}{\lambda}$$

# Wave Particle Duality

- Then in 1923, Louis de Broglie (Sorbonne, Paris) proposed a very interesting hypothesis:
- Because photons have wave and particle characteristics, perhaps all forms of matter have both properties.
- de Broglie proposed that subatomic particles, like protons and electrons, will have wavelengths just like photons.
- de Broglie suggested that a particle of mass  $m$  and velocity  $v$  would have a wavelength of:

$$\lambda = \frac{h}{p} = \frac{h}{mv}$$

← This is known as the **de Broglie wavelength** of a particle.

# Wave Particle Duality

- de Broglie also postulated that **matter waves** have **frequencies** that can be found as:

$$f = \frac{E}{h}$$

- But de Broglie only hypothesized about matter waves
  - What kind of experiment could we design to demonstrate the wave nature of particles (like electrons)?
- We could shoot electrons through a double slit apparatus and observe if there is a resulting interference pattern (similar to the pattern we observed for light).

# Wave Particle Duality

- ④ **Davisson and Germer** (US) performed this experiment in 1927.
- ④ They scattered low energy electrons from a nickel target in a vacuum.
- ④ From this experiment they found (by accident!) interference patterns and calculated a wavelength for the electron.
- ④ This wavelength agreed with the theoretical de Broglie wavelength.
- ④ This confirmed the wave nature of electrons.
- ④ This was the first experimental confirmation of the de Broglie hypothesis.

# Examples

**A.** Calculate the de Broglie wavelength of an electron ( $m_e = 9.11 \times 10^{-31} \text{ kg}$ ) moving with a speed of  $1.00 \times 10^7 \text{ m/s}$ .

**Answer:**  $\lambda = h/(m_e v) = 0.0727 \text{ nm}$

This  $\lambda$  is close to the characteristic wavelength of x-rays.

It is also on the order of the spacing of atoms in the sodium chloride lattice.

**B.** Calculate the de Broglie wavelength for a rock of mass  $50.0 \text{ g}$  thrown with a speed of  $40 \text{ m/s}$ .

**Answer:**  $3.31 \times 10^{-34} \text{ m}$

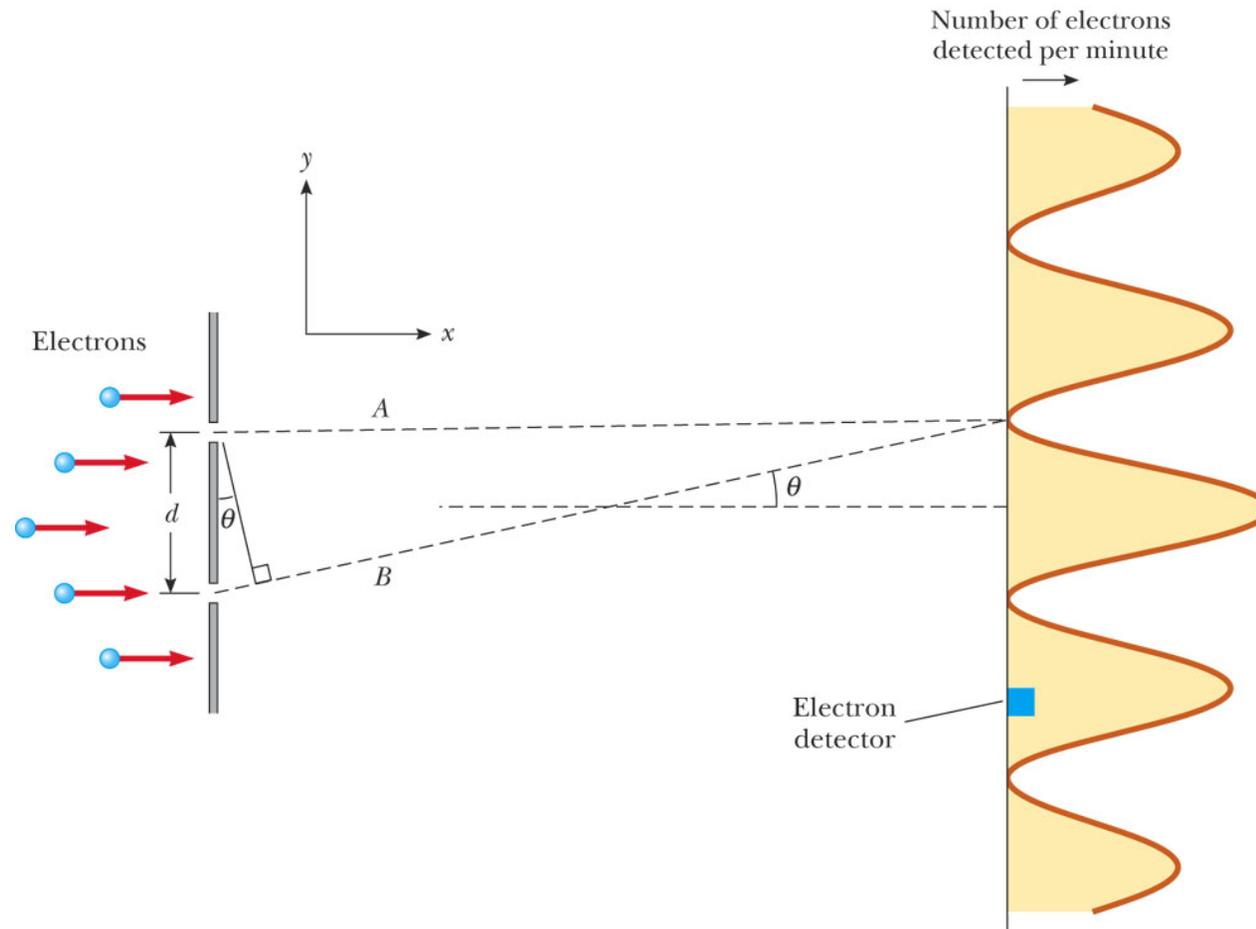
This wavelength is much smaller than the rock. Thus, the wave properties (e.g. diffraction effects) of large-scale objects cannot be observed.

# Electron Diffraction

Consider a parallel beam of monoenergetic electrons that is incident on a double slit.

Assume that the slit width are small compared to the electron wavelength.

An electron detector is positioned at a distance  $L \gg d$ .



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A typical wave interference pattern for the electron counts per minute appears.

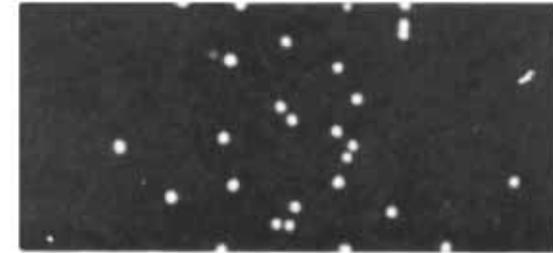
# Electron Diffraction

It is clear that electrons are interfering, which is a distinct wave-like behavior.

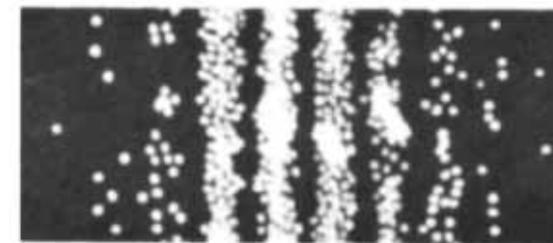
If the experiment is carried out at lower electron beam intensities, the interference pattern is still observed if the exposure is sufficiently long.

As in Chapter 27, we can use the waves in interference model to find the angular separation  $\theta$  between the central maximum and the first minimum:  $d \sin \theta = \lambda/2$ .

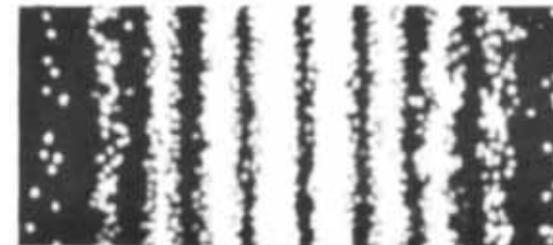
Because for electrons  $\lambda = h/p_x$  we get for small  $\theta$ :  $\sin \theta \approx \theta = h/(2p_x d)$



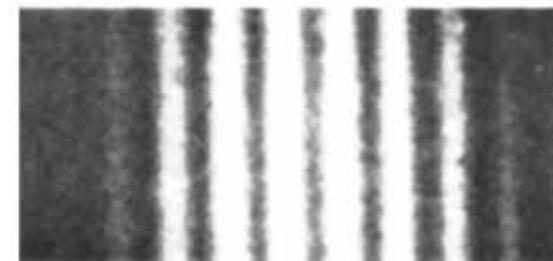
(a) After 28 electrons



(b) After 1000 electrons



(c) After 10000 electrons



(d) Two-slit electron pattern

# Electron Diffraction

- The electrons are detected as particles at a localized spot at some instance in time, but the probability of arrival at that spot is determined by finding the intensity of two interfering waves.
- We conclude that an electron interacts with both slits simultaneously.
- It is impossible to determine which slit the electron goes through.
- We can only say that the electron passes through both slits!
- If we attempt to experimentally determine which slit the electron goes through, the interference pattern is destroyed.
- The same argument applies to photons.

# The Uncertainty Principle

- ① When measurements are made, the experimenter is always faced with experimental uncertainties in the measurements.
- ② In classical mechanics, there can be measurements with arbitrarily small uncertainties (no limit).
- ③ Yet quantum mechanics predicts that a barrier to measurements with ultimately small uncertainties does exist.
- ④ When taking the measurement of position,  $x$ , the uncertainty of the measurement is given by  $\Delta x$ .
- ⑤ When taking the measurement of momentum,  $p_x$ , the uncertainty of the measurement is given by  $\Delta p_x$ .

# The Uncertainty Principle

- ④ In 1927, Werner Heisenberg unveiled his uncertainty principle:
- ④ The product of the uncertainty of the position of a particle ( $\Delta x$ ) and the uncertainty of the linear momentum of the particle ( $\Delta p_x$ ) can never be smaller than  $h/4\pi$ .
- ④ Mathematically, this becomes: 
$$\Delta x \Delta p_x \geq \frac{h}{4\pi}$$
- ④ Basically, the uncertainty principle states that **it is physically impossible to simultaneously measure the exact position and the exact linear momentum of a particle.**

# The Uncertainty Principle

- The uncertainty principle can be extended to deal with energy uncertainties as well.
- In this case, position and linear momentum are replaced with energy and time such that:

$$\Delta E \Delta t \geq \frac{h}{4\pi}$$

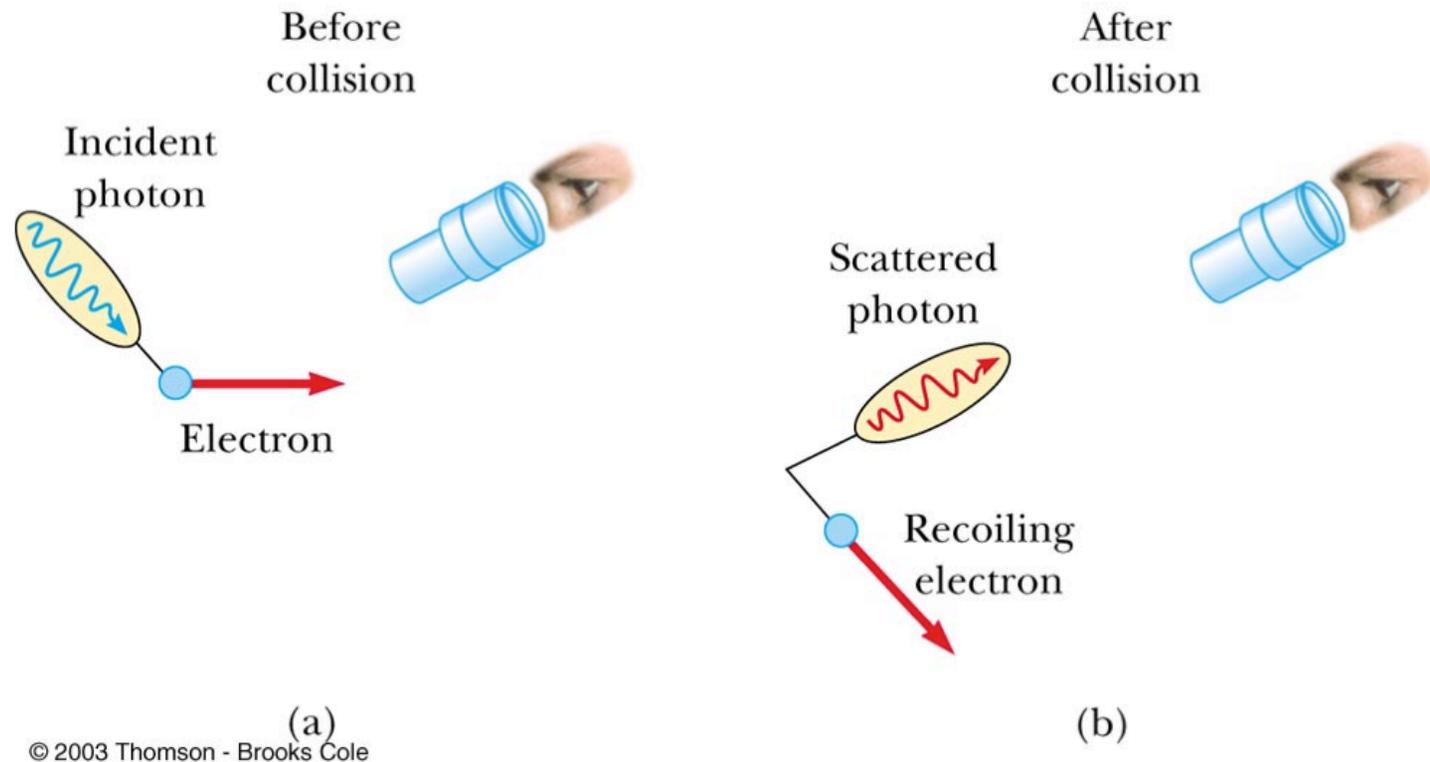
- In this form, the uncertainty principle states that it is physically impossible to measure the exact energy of a particle during a finite period of time.

# The Uncertainty Principle

One way to consider the implications of the uncertainty principle is to think about how to measure the position and linear momentum of an electron with a very powerful microscope.

In order to locate the electron accurately, at least one photon must bounce off of it.

During this interaction, momentum is transferred from the photon to the electron.



# The Uncertainty Principle

Therefore, the light that allows you to accurately locate the electron changes the momentum of the electron (the maximum change is  $\Delta p_x = h/\lambda$ ).

To minimize the momentum change, we could use longer wavelength photons.

But because the photon also has wave properties, we can determine the electron position only within one wavelength of the photon,  $\Delta x = \lambda$ .

Consequently, the position and the momentum of the electron cannot both be known precisely at the same time:  
 $\Delta p_x \Delta x \geq h$ .

Apart from the numerical factor  $1/4\pi$ , this formula is the same as Heisenberg's more precise result.

# For Next Time (FNT)

- ④ Finish reading Chapter 28
- ④ Start working on the homework for Chapter 28